Splitting methods with real-complex coefficients for separable non-autonomous semi-linear reaction-diffusion equation of **Fisher**

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Non-autonomous separable PDE

$$\frac{du}{dt} = \hat{A}(t, u) + \hat{B}(t, u), \qquad u(0) = u_0$$

- $\Rightarrow u(x,t) \in \mathbb{R}^{D}$.
- \Rightarrow (Possibly unbounded) Operators \hat{A} , \hat{B} and $\hat{A}+\hat{B}$ generate C^0 semi-groups for positive t over a finite or infinite Banach space.

Inhomogeneous Non-Autonomous Heat Equation

$$\frac{\partial u}{\partial t} = \alpha(t)\Delta u + V(x,t)u, \quad \text{or} \quad \frac{\partial u}{\partial t} = \nabla(a(x,t)\nabla u) + V(x,t)u$$

Reaction-diffusion equations

$$\frac{\partial u}{\partial t} = D(t)\Delta u + \hat{B}(t, u), \quad t \ge 0, x \in \mathbb{R}^d \quad \text{or} \quad x \in \mathbb{T}^d$$

The Problem

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References

For simplicity,

(Apparently) Linear form

$$\frac{du}{dt} = A(t)u + B(t)u,$$

 \Rightarrow A, B are the Lie operators associated to \hat{A} , \hat{B} , i.e.

$$A(t) \equiv \hat{A}(t, u) \frac{\partial}{\partial u}, \qquad B(t) \equiv \hat{B}(t, u) \frac{\partial}{\partial u}$$

which act on functions of u.

Simplest Methods For Autonomous Problems

⇒ Autonomous problem

$$\frac{du}{dt} = Au + Bu,$$

⇒ Subproblems

$$\frac{du}{dt} = Au$$
 and $\frac{du}{dt} = Bu$

 \Rightarrow Denote by e^{hA} , e^{hB} the exact *h*-flows for each problem respectively.

$$e^{hA}e^{hB}$$
 or $e^{hB}e^{hA}$,

- $S(h) = e^{h(A+B)} + O(h^2)$.
- ⇒ Strang splitting

$$S(h) = e^{h/2A} e^{hB} e^{h/2A}$$
 or $S(h) = e^{h/2B} e^{hA} e^{h/2B}$

- $S(h) = e^{h(A+B)} + O(h^3)$.
- ⇒ High-order approximations

$$\Psi(h) = e^{hb_{s+1}B} e^{ha_sA} \cdots e^{hb_2B} e^{ha_1A} e^{hb_1B},$$

at any order exist.

References

⇒ Negative coefficients for order > 2.

Splitting methods for non-autonomous problems

- A and B generate C^0 semi-groups.
- Flows e^{tA} and/or e^{tB} may not be well-defined for negative times, for instance, for the Laplacian operator

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- A and B generate C^0 semi-groups.
- Flows e^{tA} and/or e^{tB} may not be well-defined for negative times, for instance, for the Laplacian operator
- ⇒ Circumvent order-barrier, (Castella et al., BIT, 2009) and (Hansen et al., BIT, 2009) presented methods with complex coefficients having positive real parts.

 \Rightarrow Goal evaluation A(t), B(t) at real times.

Splitting method having one set of coefficients real and positive valued

$$a_i \in \mathbb{R}^+, \qquad b_i \in \mathbb{C}^+, \qquad ext{(or} \quad a_i \in \mathbb{C}^+, \qquad b_i \in \mathbb{R}^+).$$

- In order to numerically solve the system.
- ⇒ Take time as a new coordinate and split the system as (Blanes et al., JCAM, 2010)

$$\begin{cases} \frac{du}{dt}(t) = A(t_1)u \\ \frac{dt_1}{dt}(t) = 1 \end{cases} \text{ and } \begin{cases} \frac{du}{dt}(t) = B(t_1)u \\ \frac{dt_1}{dt}(t) = 0. \end{cases}$$

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Splitting methods for non-autonomous perturbed systems

Suppose that

$$||B(t)|| \ll ||A(t)||, \forall t.$$

- \Rightarrow $B = \varepsilon B_{\varepsilon}$ with $|\varepsilon| \ll 1$
- \Rightarrow A and B are qualitatively different for perturbed problems
- ⇒s-stage symmetric *BAB* compositions

$$\Psi(h) = e^{hb_{s+1}\varepsilon B} e^{ha_s A} \cdots e^{hb_2\varepsilon B} e^{ha_1 A} e^{hb_1\varepsilon B},$$

with $a_{s+1-i} = a_i$, $b_{s+2-1} = b_i$, i = 1, 2, ...

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- \Rightarrow s-stage symmetric ABA compositions

$$\Psi(h) = e^{ha_{s+1}A} e^{hb_s \varepsilon B} \cdots e^{ha_2A} e^{hb_1 \varepsilon B} e^{ha_1A},$$

with $a_{s+2-i} = a_i$, $b_{s+1-1} = b_i$, i = 1, 2, ...

 \Rightarrow For consistent methods $(\sum_i a_i = \sum_i b_i = 1)$

$$\Psi(h) = e^{h(A+\varepsilon B)+E(h,\varepsilon)}$$

$$E(h,\varepsilon) = h^3 \left(\varepsilon \, p_{aba}[[A,B],A] + \varepsilon^2 \, p_{abb}[[A,B],B]\right) + h^5 \left(\varepsilon \, p_{abaaa}[[[[A,B],A],A],A] + \mathcal{O}(\varepsilon^2)\right) + \mathcal{O}(\varepsilon h^7).$$

$$p_{aba} = rac{1}{2} \sum_{i=1}^{s} b_i c_i (1 - c_i) - rac{1}{12},$$
 $p_{abb} = \sum_{i=1}^{s} rac{1}{2} b_i^2 c_i + \sum_{1 \le i < j \le s} b_i b_j c_j - rac{1}{3},$
 $p_{abaaa} = \sum_{i=1}^{s} b_i c_i^4 - rac{1}{5}.$

$$c_i = \sum_{i=0}^i a_i,$$

$$a_0 = 0 \text{ and } c_{s+1}$$

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 and c_{s+1}

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$$i = 1, 2, ..., s$$
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$$p_{aba}=0, \qquad p_{abb}=0$$

- Fails a_i , b_i real and positive
- Goal $a_i \in \mathbb{R}^+$ and $b_i \in \mathbb{C}^+$.
- Fix $a_i \in (0,1)$ and leave the b_i to solve the order conditions.
- ⇒ It is not usually possible to concatenate the last map in one step with the first one in the following step in case of *ABA* composition.
- ⇒ In practice a *ABA* composition with the same number of exponentials as a *BAB* composition can be computationally more costly up to one additional stage.

For all these reasons we mainly focus on the BAB compositions.

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New Fourth-order methods

⇒ Fourth-order methods with a 4-stage composition

$$(b_1 a_1 b_2 a_2 b_3 a_2 b_2 a_1 b_1)$$

consistency conditions

$$a_1 + a_2 = \frac{1}{2},$$
 $2(b_1 + b_2) + b_3 = 1.$

- Fix $a_i \in (0, 1/2)$.
- Choice $a_1 = \frac{1}{4}$ obtained in (Castella et al., BIT, 2009).
- \Rightarrow Take a_1 as a free parameter to minimise the dominant error term¹

$$\min_{a_1 \in (0,1/2)} |Re(p_{abaaa})| = \min_{a_1 \in (0,1/2)} \left| \sum_{i=1}^4 Re(b_i) \, c_i^4 - \frac{1}{5} \right|.$$

 $^{^{1}}u_{n+1} = Re(\Psi(h)u_{n}).$

⇒ A 6-stage composition

$$(b_1 a_1 b_2 a_2 b_3 a_3 b_4 a_3 b_3 a_2 b_2 a_1 b_1)$$

 \Rightarrow The coefficients b_i are used to satisfy the conditions

$$p_{aba}=0, \qquad p_{abb}=0, \qquad p_{abaaa}=0.$$

- \Rightarrow We just take $a_1 = a_2 = a_3 = \frac{1}{6}$.
- → Methods which derived in this work can be seen our paper (Blanes et al., Submitted).

Splitting methods

⇒ Standard Strang decomposition for non-autonomous system

$$S(h) = e^{h/2 B_1} e^{hA_0} e^{h/2 B_0}$$
 (1)

$$\Rightarrow$$
 $A_i = A(t_n + ih), B_i = B(t_n + ih).$

Split the system

$$\begin{cases} \frac{du}{dt}(t) = A(t_1)u \\ \frac{dt_1}{dt}(t) = 0 \\ \frac{dt_2}{dt}(t) = 1 \end{cases} \quad \text{and} \quad \begin{cases} \frac{du}{dt}(t) = B(t_2)u \\ \frac{dt_1}{dt}(t) = 1 \\ \frac{dt_2}{dt}(t) = 0. \end{cases}$$

⇒ Fourth order extrapolation method

$$\Phi^{[4]}(h) = \frac{4}{3} e^{\frac{h}{4}B_1} e^{\frac{h}{2}A_{1/2}} e^{\frac{h}{2}B_{1/2}} e^{\frac{h}{2}A_0} e^{\frac{h}{4}B_0} - \frac{1}{3} e^{\frac{h}{2}B_1} e^{hA_0} e^{\frac{h}{2}B_0}. \quad (2)$$

- ⇒ Following schemes are then considered:
 - ⇒ with real coefficients
 - Strang: The second-order symmetric Strang splitting method (1);
 - **(6,2)**: The symmetric splitting method of effective order (6,2) given in (McLachlan, BIT, 1995);
 - (EXT4): The fourth-order extrapolation method (2);
 - \Rightarrow with complex coefficients and $a_i \in \mathbb{R}^+$
 - (RC4): The 4-stage fourth-order method from (Castella et al., BIT, 2009);
 - (O4): The 4-stage fourth-order method built in (Blanes et al., MATCMP, 2013), whose coefficients are available at http://www.gicas.uji.es/Research/splitting-complex.html and referred as "Order 4 (optimized)";
 - **(SM4)**: The new optimized 4-stage fourth-order method given in (Blanes et al., Submitted);
 - (SM(6,4)): The new 6-stage fourth-order method are given in (Blanes et al., Submitted);

The semi-linear reaction-diffusion equation of Fisher

$$\frac{\partial u}{\partial t} = \alpha(t)^2 \Delta u + F(u,t), \qquad u(x,0) = u_0(x),$$

- \Rightarrow Periodic boundary conditions in the space domain [0, 1].
- ⇒ Fisher's potential

$$F(u) = \gamma(t)u(1-u), \qquad \gamma(t) = \frac{2-e^{-\beta t}}{100}$$

- $\Rightarrow \alpha(t) = \frac{1}{4} + \mu \cos(wt).$
- → Discretization in space,

$$\frac{dU}{dt} = \alpha(t)^2 AU + F(U, t),$$

- $\Rightarrow U = (U_1, \ldots, U_N) = (u_1, \ldots, u_N) \in \mathbb{R}^N.$
- \Rightarrow A is $N \times N$ circulant matrix.
- $\Rightarrow F(U,t) = \gamma(t)(U_1(1-U_1),...,U_N(1-U_N)).$

Splitting technique

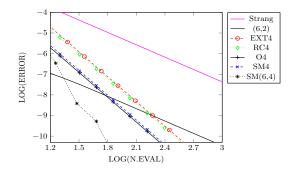
$$\begin{cases} \frac{dU}{dt}(t) = \alpha(t_1)^2 A \Delta U \\ \frac{dt_1}{dt}(t) = 1 \end{cases} \text{ and } \begin{cases} \frac{dU}{dt}(t) = \gamma_1 U(1 - U) \\ \frac{dt_1}{dt}(t) = 0. \end{cases}$$

$$\Rightarrow \gamma_1 = \gamma(t_1)$$

Analytic solution of scalar equations

$$u(x,h) = u_0(x) + u_0(x)(1 - u_0(x)) \frac{(e^{\gamma_1 h} - 1)}{1 + u_0(x)(e^{\gamma_1 h} - 1)}$$

 $\Rightarrow u(x, h)$ is well-defined for small complex time h.



$$\Rightarrow u_0(x) = \sin(2\pi x).$$

$$\Rightarrow \beta = 1, \mu = 1/8, w = 4, N = 100$$
 at final time $t = 1$.

 \Rightarrow Compute the error (in the 2-norm) at the final time t=1 by applying the same composition methods as in the linear case.

 \Rightarrow N.EVAL is number of evaluations of the flow A.

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