# Geometric integrators with complex coefficients

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# Introduction

The numerical integration of non-reversible systems using high order splitting methods with complex coefficients (having positive real part) has been recently considered Example: The linear heat equation with potential

$$\frac{\partial}{\partial t}\mathbf{u} = \triangle \mathbf{u} + V(x)\mathbf{u}$$

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Some questions:

- ∃ methods at any order with coefficients having positive real parts?
- Are the new methods efficient?
- Are useful for long time integrations? Backward error analysis?

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- From the theoretical point of view:
  - Singularities in the complex domain can appear
  - Bounded solutions in the real space are unbounded in the complex domain, and this can affect to the stability of the methods

 The evolution of the solution on the extended manifold in the complex domain needs to be studied

## Example: The Volterra-Lotka problem

Let us consider the Volterra–Lotka problem

$$\dot{u} = u(v-2), \qquad \dot{v} = v(1-u)$$

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First integral:  $I(u, v) = \ln(uv^2) - (u + v)$ .

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First integral:  $l(u, v) = \ln(uv^2) - (u + v)$ . It can be considered as a Hamiltonian system with

$$H = (2p - e^p) + (q - e^q)$$

with  $q = \ln u$ ,  $p = \ln v$ Split:  $f_A = (u(v-2), 0)$ ,  $f_B = (0, v(1-u))$ 

Given a symmetric 2nd order  $S^{[2]}$  one gets a 4th order integrator  $S^{[4]} : \mathbb{R}^{2d} \to \mathbb{R}^{2d}$  as

$$S_h^{[4]} = S_{\alpha h}^{[2]} \circ S_{\beta h}^{[2]} \circ S_{\alpha h}^{[2]},$$
$$2\alpha + \beta = 1, \qquad 2\alpha^3 + \beta^3 = 0$$

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$$\alpha = \frac{1}{2 - 2^{1/3}} \simeq 1.35 > 1, \qquad \beta = 1 - 2\alpha \simeq -1.7 < 0$$

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 $\alpha \simeq 0.324 \pm i0.135, \qquad \beta \simeq 0.351 \mp i0.269$ 

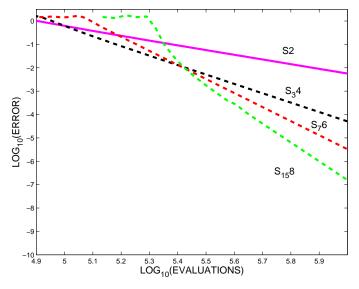
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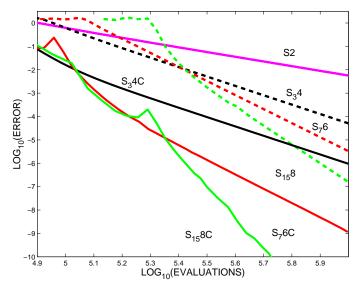
 $lpha \simeq 0.324 \pm i0.135, \qquad eta \simeq 0.351 \mp i0.269$ Initial conditions ( $u_0, v_0$ ) = (4,2),  $t \in [0, 1000]$ 

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Order conditions

$$\sum_{i=1}^k \alpha_i = 1, \qquad \sum_{i=1}^k \alpha_i^3 = 0$$

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If  $\alpha_i \in \mathbb{C}$  then

$$\max_{i=1,\ldots,k} \operatorname{Arg}(\alpha_i) - \min_{i=1,\ldots,k} \operatorname{Arg}(\alpha_i) \geq \frac{\pi}{3}.$$

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Starting from  $\mathcal{S}^{[2]}(h)$ , we have

$$\mathcal{S}^{[2(p+1)]}(h) = \prod_{i_p=1}^{m_p} \left( \prod_{i_{p-1}=1}^{m_{p-1}} \left( \dots \left( \prod_{i_1=1}^{m_1} \mathcal{S}^{[2]}(\alpha_{p,i_p} \alpha_{p-1,i_{p-1}} \cdots \alpha_{1,i_1} h) \right) \dots \right) \right)$$

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Castella, Chartier, Descombes, & Vilmart, BIT 49 (2009), 487-508, and Hansen & Ostermann, BIT 49 (2009), 527-542, obtained methods up to order 14 with coefs. having positive real part.

$$\mathcal{S}_{h}^{[2]} \to \mathcal{S}_{h}^{[4]} \to \mathcal{S}_{h}^{[6]} \to \mathcal{S}_{h}^{[8]} \to \dots \mathcal{S}_{h}^{[14]} \to \mathcal{S}_{h}^{[16]}$$

#### Lemma

For  $k \ge 2$  and  $r \ge 2$ , consider  $(z_1, \ldots, z_k) \in (\mathbb{C}_+)^k$  such that  $\sum_{i=1}^k z_i^r = 0$ . Then we have

$$\max_{i=1,\ldots,k} \operatorname{Arg}(z_i) - \min_{i=1,\ldots,k} \operatorname{Arg}(z_i) \geq \frac{\pi}{r}.$$

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#### Proof.

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$$\max_{i=1,\ldots,k} \operatorname{Arg}(z_i) - \min_{i=1,\ldots,k} \operatorname{Arg}(z_i) < \frac{\pi}{r}.$$

then

$$\max_{i=1,\ldots,k} \operatorname{Arg}(Z_i^r) - \min_{i=1,\ldots,k} \operatorname{Arg}(Z_i^r) < \pi.$$

and obviously  $\sum_{i=1}^{k} z_i^r \neq 0$ .

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#### Theorem

Starting from a second-order symmetric method  $S^{[2]}(h)$ , all methods  $S^{[2p]}(h)$  of order 2p = 16, 18, ... from the previous recursion have at least one coefficient with negative real part.

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#### Proof.

We assume that all methods  $S^{[2q]}(h)$ , q = 1, ..., p have all their coefficients in  $\mathbb{C}_+$ . Using Lemma 1 we have

$$\forall q = 1, \dots, p, \quad \max_{i=1,\dots,m_q} \operatorname{Arg}(\alpha_{q,i}) - \min_{i=1,\dots,m_q} \operatorname{Arg}(\alpha_{q,i}) \geq \frac{\pi}{2q+1},$$

#### so that

$$\max_{i_1,\dots,i_p} \operatorname{Arg}\left(\prod_{j=1}^p \alpha_{j,i_j}\right) - \min_{i_1,\dots,i_p} \operatorname{Arg}\left(\prod_{j=1}^p \alpha_{j,i_j}\right) \ge \frac{\pi}{3} + \dots + \frac{\pi}{2p+1}.$$
  
Since  $\frac{1}{3} + \frac{1}{5} + \dots + \frac{1}{15} > 1$ , then  $2p = 14$  is an upper bound.  $\Box$ 

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so that

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Since  $\frac{1}{3} + \frac{1}{5} + \dots + \frac{1}{15} > 1$ , then  $2p = 14$  is an upper bound.  $\Box$ 

This is a sharp bound since methods of order 14 have been obtained

From the computational point of view, it is more efficient to build methods directly by the composition

$$\mathcal{S}^{[2p]}(h) = \mathcal{S}^{[2]}(\gamma_1 h) \cdots \mathcal{S}^{[2]}(\gamma_s h)$$

We have built methods of order 6 and 8:

$$S^{[6]}(h) = S^{[2]}(\gamma_1 h) \cdots S^{[2]}(\gamma_7 h) S^{[8]}(h) = S^{[2]}(\gamma_1 h) \cdots S^{[2]}(\gamma_{15} h)$$

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We have also built methods of order 16 with coefficients having their real part positive. The procedure followed is:

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∃ methods at all orders? We are still ignorant, but at a higher level of ignorance!

## Example: The Volterra-Lotka problem

$$\dot{u} = u(v-2), \qquad \dot{v} = v(1-u)$$

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Measure the relative error:  $|I - I_0|/|I_0|$  with  
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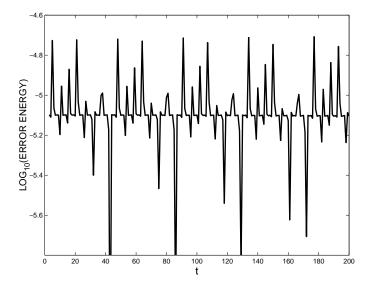
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Initial conditions  $(u_0, v_0) = (4, 2)$ , time step:  $h = \frac{1}{8}$ Measure the relative error:  $|I - I_0| / |I_0|$  with  $I(u, v) = \ln(uv^2) - (u + v)$ .

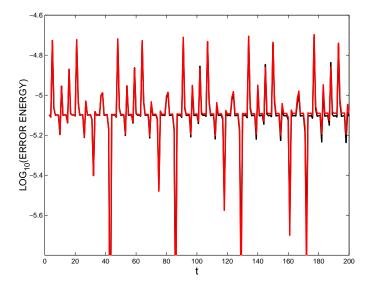
- I) Projection at the end of the integration
- II) Projection at each step

# Example:

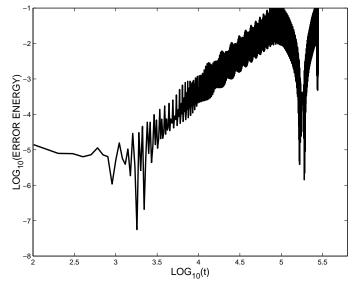


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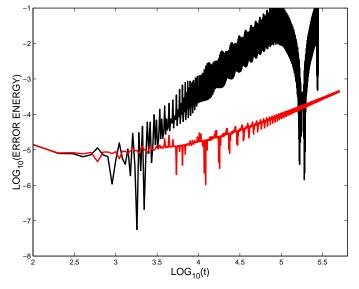
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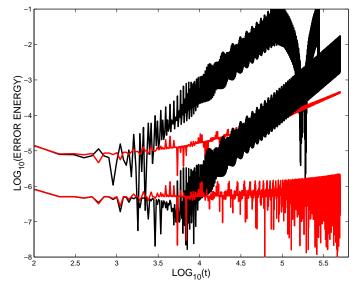
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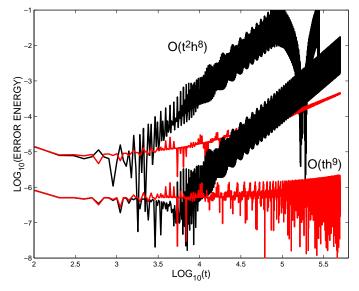
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Let  $\Phi(h) = e^{hF}$  denote the exact solution and  $S_r(h)$  a with complex coefficients method of order  $r = \min\{q, p\}$  such that

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$$S_r(h) = \exp\left(hF + h^{q+1}F_R + ih^{p+1}F_I\right)$$

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 $F_R$ ,  $F_I$  are elements of the Lie algebra associated to the components of F, but  $\hat{F}_I^2$  is not in the Lie algebra.

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$$Re\left(S_{\rho}^{N}(h)\right) = \left(I - t^{2}h^{2\rho}\check{F}_{I}^{2} + \ldots\right)\exp\left(t(F + h^{q}\check{F}_{R})\right)$$

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while to project at each step corresponds to

$$Re\left(S_{p}(h)\right)^{N} = \left(I - th^{2p+1}\check{F}_{I}^{2} + \ldots\right) \exp\left(t(F + h^{q}\check{F}_{R})\right)$$

#### Some other problems of interest

(a) The linear Schrödinger equation ( $\hbar = 1$ ):

$$\begin{split} i\frac{\partial}{\partial t}\Psi(x,t) &= \left(-\frac{1}{2m}\nabla^2 + V(x)\right)\Psi(x,t)\\ \mathbf{u}(h) &= e^{ih(\Delta+\mathbf{V})}\mathbf{u}_0\\ \text{Methods with} \quad \frac{\mathbf{a}_i \in \mathbb{R}^+, \ \mathbf{b}_i \in \mathbb{C}. \end{split}$$

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Methods with  $a_i \in \mathbb{R}^+, b_i \in \mathbb{C}$ .

(b) The LSE integrated in the pure imaginary time

$$\frac{\partial}{\partial \tau} \Psi(x,\tau) = \left(\frac{1}{2m} \nabla^2 - V(x)\right) \Psi(x,\tau)$$

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Methods with  $a_i$ ,  $b_i \in \mathbb{C}_+$ .

In addition  $[V, [V, [V, \nabla^2]]] = 0$ 

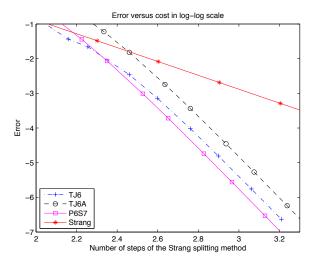
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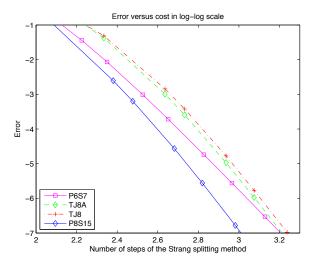
#### Let us consider the scalar equation in one-dimension

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$$\frac{\partial u}{\partial t} = \Delta u + (2 + \sin(2\pi x)) u$$
$$u(x,0) = \sin(2\pi x)$$
$$x \in [0,1] \text{ with } N = 100$$
periodic boundary conditions.



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