

# Symplectic methods for the time integration of the Schrödinger equation

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Joint work with:

[Fernando Casas](#) and [Ander Murua](#)

# The Goal

The numerical integration of the **time-dependent Schrödinger Equation** ( $\hbar = 1$ )

$$i\frac{\partial}{\partial t}\psi(x, t) = -\frac{1}{2\mu}\nabla^2\psi(x, t) + V(x)\psi(x, t)$$

$$\psi(x, 0) = \psi_0(x), \quad x \in \mathbb{R}^D, \quad t \in [0, T].$$

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We consider the case where an appropriate spatial discretisation is used and one has to solve an **IVP**

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# The Numerical Integration of Differential Equations

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About 50 years ago, researchers had the hope to develop a few numerical methods to cover the numerical solution of most problems, i.e. to build a **black box** with a few number of methods implemented.

Soon, it was clear that this was too **optimistic** due to the huge variety of problems of very different nature, and started to look for **methods tailored** for different classes of **problems**

# Different families of methods

- Runge–Kutta methods (explicit and implicit)
- Multistep methods (explicit and implicit)
- Extrapolation methods
- etc.



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However, most equations to be solved originate from physical problems obtained from **First Principles**, which makes the solutions to have very particular qualitative properties.

## Geometric Integration

- **Symplectic Integrators**
- Lie group methods
- Volume-preserving methods
- etc.

# The numerical Solution of very particular problems

The development of computers allowed to researchers in physics, chemistry, engineering, etc. to study more challenging problems from the **computational point of view**.

These problems can not be solved by the computer just by brute force, and **tailored methods** have to be developed.

# Example

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- for **60 Myrs**
- **Backward in time**
- to very **high accuracy**

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Actual methods (already tailored for this problem) allowed for a faithful integration over **40 Myrs**, with good agreement with observations by geologists.

We were asked to develop new methods with better performance than the existing ones for this particular problem.

# Example:

		Motivation	Intro. Symp.	Specific Symp.	Test										
Eonothem Eon	Erathem Era	System Period	Series Epoch	Stage Age	Age Ma	GSSP	Eonothem Eon	Erathem Era	System Period	Series Epoch	Stage Age	Age Ma	GSSP	Eonothem Eon	Erathem Era
Phanerozoic	Cenozoic	Neogene	Holocene					Phanerozoic	Mesozoic	Upper	Tithonian	145.5 ± 4.0		Phanerozoic	Paleozoic
				Upper	0.0 - 0.15		Kimmeridgian				150.8 ± 4.0				
			Pleistocene	Middle	0.126		Oxfordian				155.7 ± 4.0				
				Lower	0.781	👉	Callovian				161.2 ± 4.0				
					1.806	👉	Barremian				164.7 ± 4.0				
			Pliocene	Gelasian	2.588	👉									
		Piacenzian		3.600	👉										
		Miocene	Zanclean	5.332	👉										
			Messinian	7.246	👉										
			Tortonian	11.608	👉										
	Serravallian		13.65		Santonian	196.5 ± 1.0	👉								
	Langhian		15.97		Hettangian	199.6 ± 0.6									
	Burdigalian		20.43		Rhaetian	203.6 ± 1.5									
	Paleogene	Oligocene	Aquitanian	25.03	👉	Norian	216.5 ± 2.0								
			Chatthian	28.4 ± 0.1	👉	Carnian	228.0 ± 2.0								
			Rupelian	33.9 ± 0.1	👉	Ladinian	237.0 ± 2.0								
			Priabonian	37.2 ± 0.1											
		Eocene	Bartonian	40.4 ± 0.2											
			Lutetian	48.6 ± 0.2	👉										
			Ypresian	55.8 ± 0.2	👉										
Paleocene		Thanetian	58.7 ± 0.2												
		Selandian	61.7 ± 0.2	👉	Capitanian	260.0 ± 2.0	👉								
		Danian	65.5 ± 0.3	👉											

**Astronomical calibration (2004)**  
(La2004 : 40 Ma)

**Astronomical calibration (in project)**

GTS2004 : Gradstein, Ogg, Smith, 2004

**B**, Casas, Farrés, Makazaga, Murua and Laskar. Two submitted papers.

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We present with some detail the steps to follow in order to look for efficient methods for some problems in Quantum Mechanics.

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- 4 (**Practical**) To collaborate with experts on these fields

# Back to the Physical Problem

We illustrate this procedure on the 1-dim SE

$$i \frac{\partial}{\partial t} \psi(x, t) = \left( -\frac{1}{2\mu} \frac{\partial^2}{\partial x^2} + V(x) \right) \psi(x) = \mathcal{H} \psi(x, t)$$

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$\psi(x, 0) = \psi_0(x)$ .  $\mathcal{H}$  is an **Hermitian operator**. Then

$$\mathcal{H} \varphi_k(x) = \mathcal{E}_k \varphi_k(x), \quad k = 0, 1, 2, \dots$$

where  $\{\mathcal{E}_k, \varphi_k(x)\}$  are the real eigenvalues and orthonormal eigenfunctions, and

$$\psi_0(x) = \sum_{k=0}^{\infty} c_k \varphi_k(x) \quad \Rightarrow \quad \psi(x, t) = \sum_{k=0}^{\infty} c_k e^{-it\mathcal{E}_k} \varphi_k(x)$$

# Back to the problem

- $|\psi(x, t)|^2$ : probability to find the quantum particle in  $(x, t)$
- Then,  $\psi(x, t) \rightarrow 0$  as  $x \rightarrow \pm\infty$
- It suffices to consider a bounded region where the solution and all its derivatives vanishes at the boundaries (**periodic problem**)
- We can use spectral methods for the spatial discretisation

# Back to the problem

A mesh with  $d$  points  $\Rightarrow$   $d$ -dimensional linear problem

$$i \frac{d}{dt} u = H u \quad \Rightarrow \quad u(T) = e^{-i T H} u(0)$$

where  $u \in \mathbb{C}^d$  and  $H \in \mathbb{R}^{d \times d}$  is a **Hermitian matrix**.

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$$H v_k = E_k v_k, \quad k = 0, 1, 2, \dots, d-1$$

where we expect

$$E_k \simeq \mathcal{E}_k, \quad v_{k,j} \simeq \psi_k(x_j)$$

for  $k = 0, 1, \dots, d_0 - 1$  with  $d_0 \leq d$ , and

$$u_0 = \sum_{k=0}^{d_0-1} \hat{c}_k v_k, \quad \Rightarrow \quad u(T) = \sum_{k=0}^{d_0-1} \hat{c}_k e^{-iT E_k} v_k$$



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We look for **polynomial approximations** to the exponential. Formally, the problem to solve is

$$i \frac{du}{dt} = P^{-1} \begin{pmatrix} E_0 & & & \\ & E_1 & & \\ & & \ddots & \\ & & & E_{d-1} \end{pmatrix} P u = H u$$

which is just a set of  **$d$  harmonic oscillators**

**The goal:** To approximate the solution

$$u(T) = e^{-iTH} u_0$$

using an algorithm which involves vector-matrix products.

Inputs:

- $T$ : Time of integration.
- $u_0$ : Initial conditions.
- $Hu$ : A subroutine, to compute the product of a vector  $u$  with the matrix  $H$ .
- $E_{min}, E_{max}$ , such that

$$E_{min} \leq E_0 < \dots < E_{d-1} \leq E_{max}$$

(they are usually known).

- $tol$ : The approximated solution,  $\tilde{u}$ , must satisfy

$$\|u(T) - \tilde{u}\| < tol$$

We can take a shift to the center of the eigenvalues

$$e^{-itH} = e^{-iT\alpha} e^{-iT(H-\alpha I)}$$

with

$$\alpha = \frac{E_{\max} + E_{\min}}{2}$$

and a normalization

$$\exp(-iTH) = \exp(-iT\alpha) \exp(-iT\beta\tilde{H})$$

where

$$\tilde{H} = \frac{H - \alpha I}{\beta}, \quad \beta = \frac{E_{\max} - E_{\min}}{2}$$

so

$$-1 \leq \sigma(\tilde{H}) \leq 1$$

# The Mathematical Problem

To approximate

$$w(T) = e^{-iT\beta\tilde{H}}u_0$$

where

$$\beta = \frac{E_{\max} - E_{\min}}{2}, \quad -1 \leq \sigma(\tilde{H}) \leq 1$$

$$\tau = T\beta \quad (\text{effective time})$$

using polynomial approximations (vector matrix products) such that

$$\|w(T) - w_{ap}\| < tol$$

# The State of the Art: Taylor method

$$(u = q + ip)$$

$$w_T = \sum_{k=0}^m \frac{(-iT\beta)^k}{k!} \tilde{H}^k u_0 = (T_m^C - iT_m^S)(q_0 + ip_0)$$

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## Horner's algorithm

$$y_0 = u_0$$

**do**  $k = 1, m$

$$y_k = u_0 - i \frac{T\beta}{m+1-k} \tilde{H} y_{k-1}$$

**enddo**

$$w_T = y_m$$



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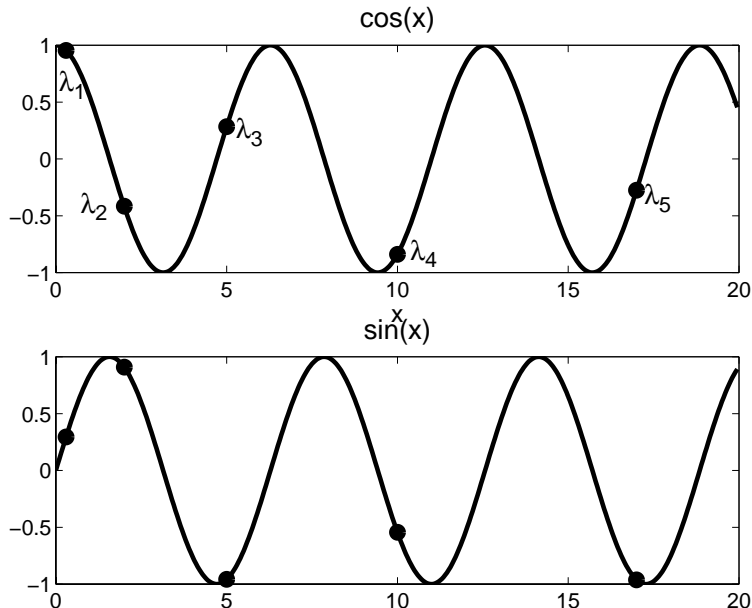
## Horner's algorithm

```
y0 = u0  
do k = 1, m  
    yk = u0 - i  $\frac{T\beta}{m+1-k}$   $\tilde{H}$  yk-1  
enddo  
wT = ym
```

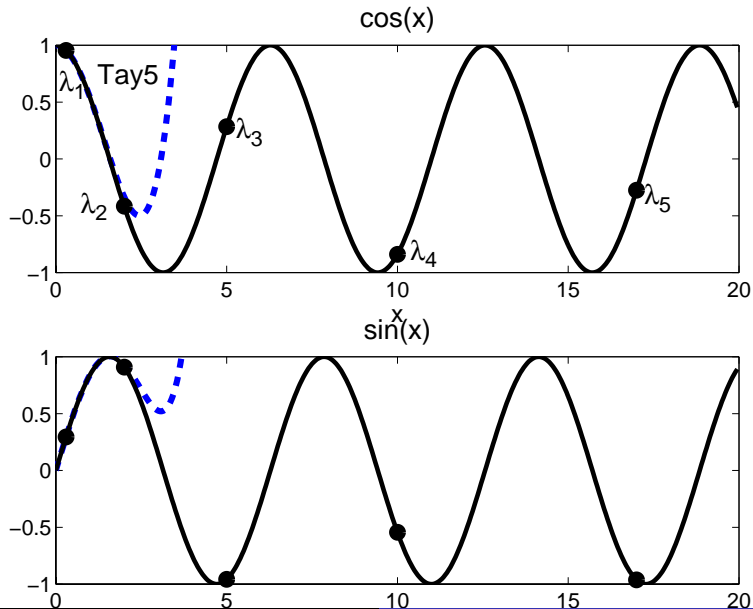
## Error bounds

$$\|w(T) - w_T\| < \frac{(T\beta)^{m+1}}{(m+1)!} e^{T\beta}, \quad \frac{T\beta}{m} \lesssim 0.3$$

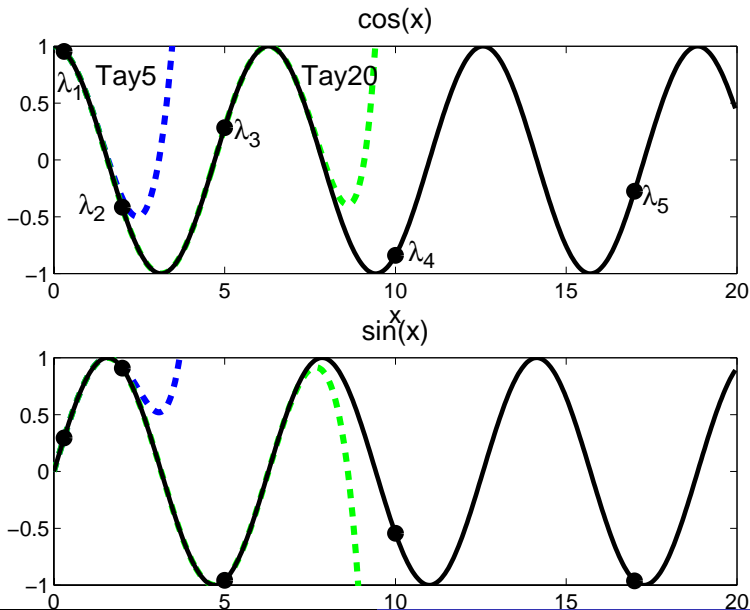
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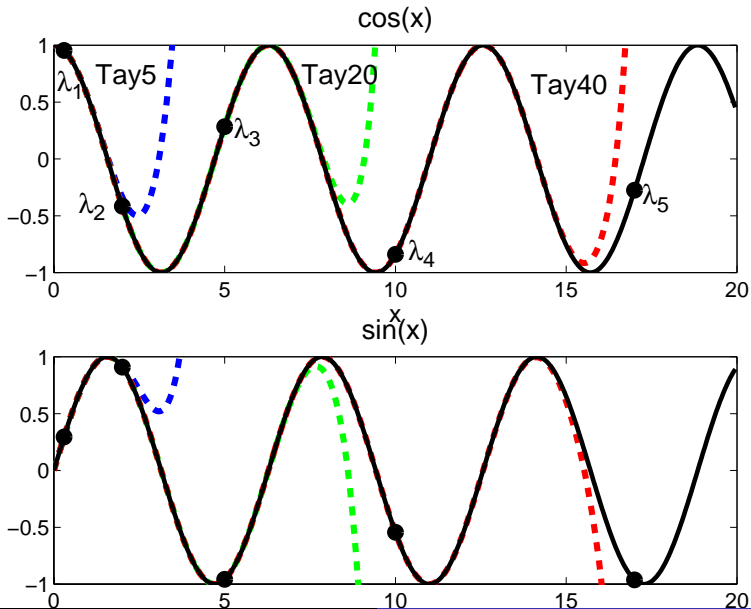
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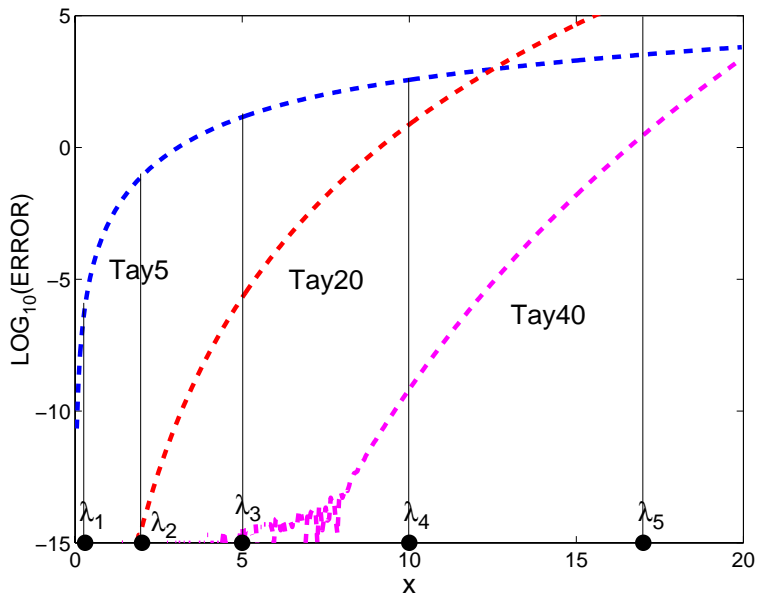
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# The State of the Art: Chebyshev method

$$w_C = \left( J_k(T\beta) + 2 \sum_{k=1}^m (-i)^k J_k(T\beta) T_k(\tilde{H}) \right) u_0 = (C_m^C - iC_m^S)(q_0 + ip_0)$$

$J_k(t)$ : Bessel functions of the first kind

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The **Clenshaw algorithm** ( $c_k = (-i)^k J_k(\tau\beta)$ ):

$$d_{m+2} = 0, \quad d_{m+1} = 0$$

**do**  $k = m, 0$

$$d_k = c_k u_0 + 2\tilde{H}d_{k+1} - d_{k+2}$$

**enddo**

$$w_C \equiv P_{m-1}^C(\tau\tilde{H}) u_0 = d_0 - d_2$$



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**Error bounds**

$$\|w(T) - w_C\| < 4 \left( e^{1 - (\beta\tau/2(m+1))^2} \frac{\beta\tau}{2(m+1)} \right)^{(m+1)}, \quad \frac{T\beta}{m} < 1$$

# The State of the Art: Symplectic methods

$$u = e^{-iT\beta\tilde{H}}u_0 \Rightarrow q + ip = (\cos(T\beta\tilde{H}) - i\sin(T\beta\tilde{H}))(q_0 + ip_0)$$

$$\begin{Bmatrix} q \\ p \end{Bmatrix} = \begin{pmatrix} \cos(T\beta\tilde{H}) & \sin(T\beta\tilde{H}) \\ -\sin(T\beta\tilde{H}) & \cos(T\beta\tilde{H}) \end{pmatrix} \begin{Bmatrix} q_0 \\ p_0 \end{Bmatrix}$$

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## Splitting Symplectic methods

$$\begin{Bmatrix} q_S \\ p_S \end{Bmatrix} = \prod_{k=1}^m \begin{pmatrix} I & 0 \\ -b_k T\beta\tilde{H} & I \end{pmatrix} \begin{pmatrix} I & a_k T\beta\tilde{H} \\ 0 & I \end{pmatrix} \begin{Bmatrix} q_0 \\ p_0 \end{Bmatrix}$$

Gray & Manolopoulos J. Chem. Phys. (1996):  $m = 2, 4, 6, 8, 10, 12$   
The [algorithm](#):

**do**  $k = 1, m$

$q := q + a_k T\beta\tilde{H}p$

$p := p - b_k T\beta\tilde{H}q$

**enddo**

$$\frac{T\beta}{m} < 2$$

NO Error bounds

# Numerical example 1

(Lubich, [Blue book](#), 2008) To approximate

$$e^{-iH}u_0$$

with  $u_0$  a unitary random vector and

$$H = \frac{\lambda}{2} \begin{pmatrix} 2 & -1 & & & & \\ -1 & 2 & -1 & & & \\ & & \ddots & & & \\ & & & -1 & 2 & -1 \\ & & & & -1 & 2 \end{pmatrix} \in \mathbb{R}^{N \times N}, \quad N = 10000$$

$0 \leq E_k \leq 2\lambda$ ,  $k = 1, 2, \dots, 10000$

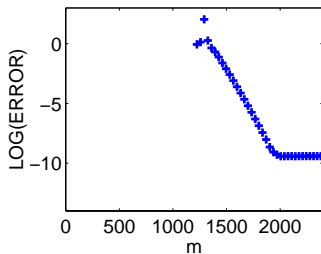
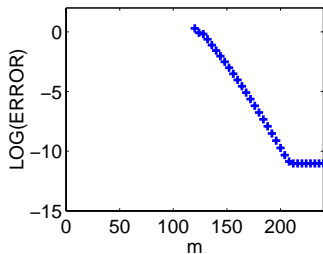
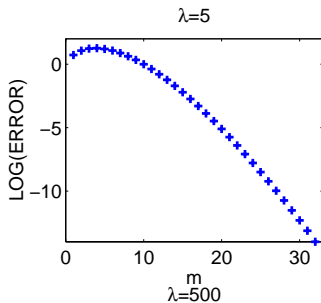
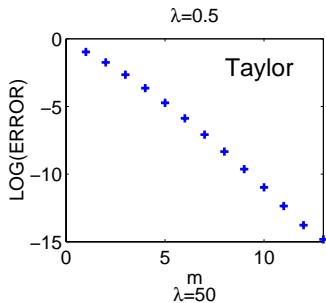
After a shift,  $H - \lambda I$ , we can take:  $T\beta = \lambda$

We approximate:

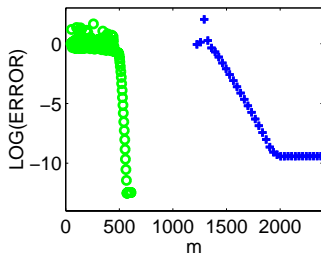
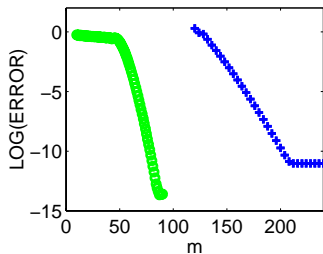
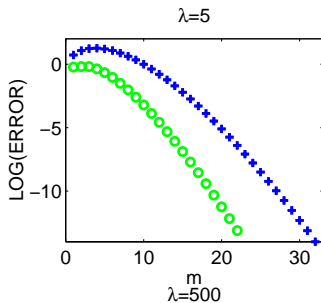
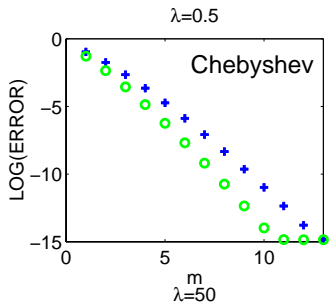
$$e^{-i\lambda} e^{-i\lambda \hat{H}} u_0, \quad \hat{H} = (H - \lambda I)/\lambda$$

$$\lambda = 0.5, \quad \lambda = 5, \quad \lambda = 50, \quad \lambda = 500$$

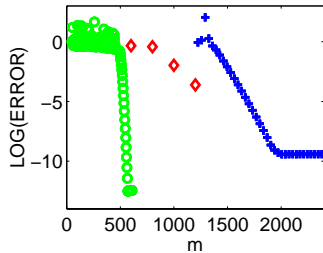
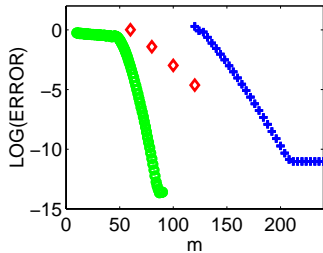
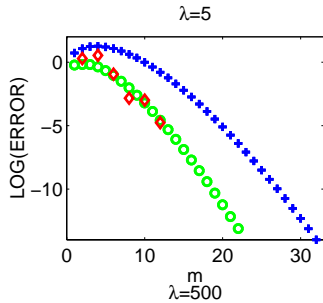
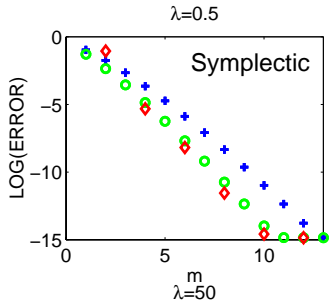
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## Numerical example 2: The scalar problem

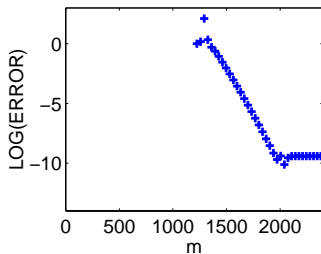
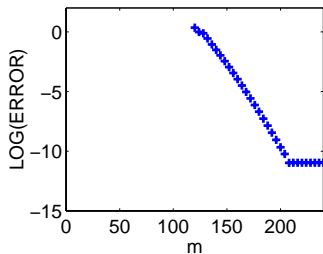
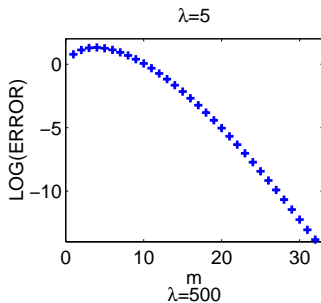
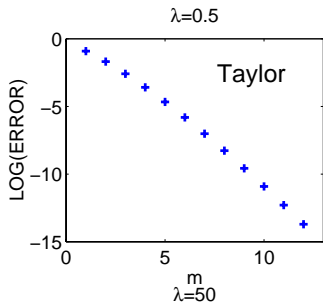
$$e^{-i\lambda} u_0, \quad \lambda \in \mathbb{R}$$

with  $u_0 = 1$

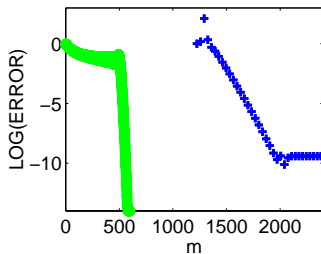
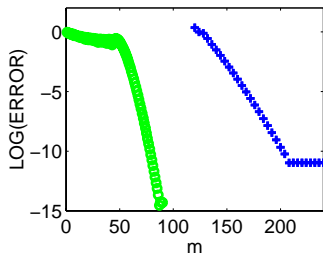
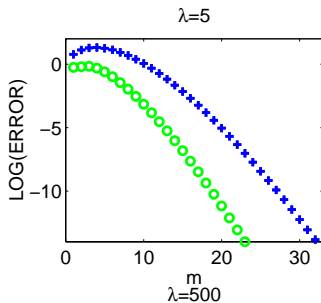
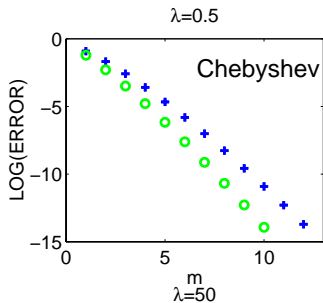
$$\lambda = 0.5, \quad \lambda = 5, \quad \lambda = 50, \quad \lambda = 500$$



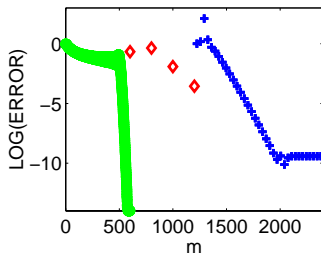
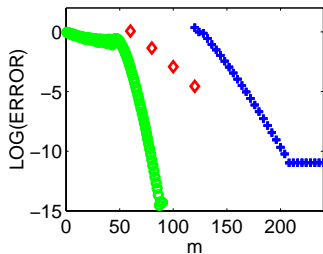
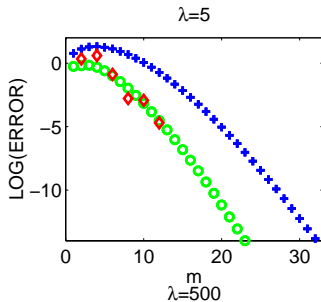
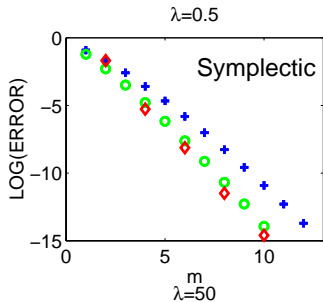
# Example 2:



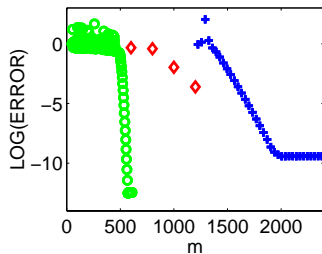
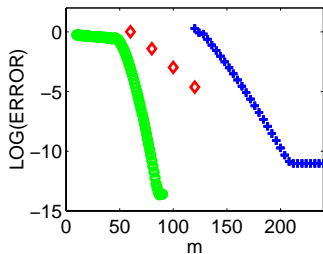
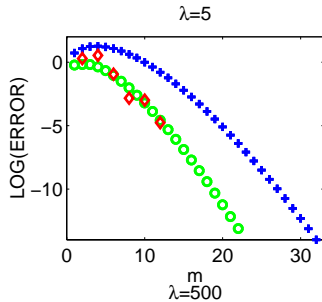
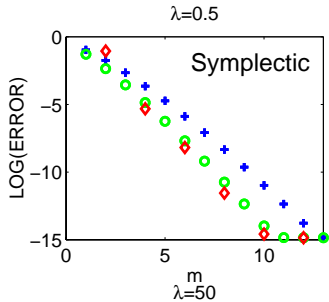
# Example 2:



# Example 2:



# Example 1:



# The Simplified Model

$$u = e^{-iy} u_0 \quad \Rightarrow \quad q + ip = (\cos(y) - i \sin(y))(q_0 + ip_0)$$

$$\begin{Bmatrix} q \\ p \end{Bmatrix} = \begin{pmatrix} \cos(y) & \sin(y) \\ -\sin(y) & \cos(y) \end{pmatrix} \begin{Bmatrix} q_0 \\ p_0 \end{Bmatrix}$$

## Taylor method

$$\begin{Bmatrix} q_T \\ p_T \end{Bmatrix} = \begin{pmatrix} T_1^m & T_2^m \\ -T_2^m & T_1^m \end{pmatrix} \begin{Bmatrix} q_0 \\ p_0 \end{Bmatrix}$$

$$\frac{T\beta}{m} < 0.3$$

## Taylor method

$$\begin{Bmatrix} q_T \\ p_T \end{Bmatrix} = \begin{pmatrix} T_1^m & T_2^m \\ -T_2^m & T_1^m \end{pmatrix} \begin{Bmatrix} q_0 \\ p_0 \end{Bmatrix}$$

$$\frac{T\beta}{m} < 0.3$$

## Chebyshev method

$$\begin{Bmatrix} q_C \\ p_C \end{Bmatrix} = \begin{pmatrix} C_1^m & C_2^m \\ -C_2^m & C_1^m \end{pmatrix} \begin{Bmatrix} q_0 \\ p_0 \end{Bmatrix}$$

$$\frac{T\beta}{m} < 1$$

## Taylor method

$$\begin{Bmatrix} q_T \\ p_T \end{Bmatrix} = \begin{pmatrix} T_1^m & T_2^m \\ -T_2^m & T_1^m \end{pmatrix} \begin{Bmatrix} q_0 \\ p_0 \end{Bmatrix} \quad \frac{T\beta}{m} < 0.3$$

## Chebyshev method

$$\begin{Bmatrix} q_C \\ p_C \end{Bmatrix} = \begin{pmatrix} C_1^m & C_2^m \\ -C_2^m & C_1^m \end{pmatrix} \begin{Bmatrix} q_0 \\ p_0 \end{Bmatrix} \quad \frac{T\beta}{m} < 1$$

## Symplectic methods

$$\begin{pmatrix} 1 & 0 \\ -b_k y & 1 \end{pmatrix} \begin{pmatrix} 1 & a_k y \\ 0 & 1 \end{pmatrix} = \begin{pmatrix} 1 & a_k y \\ -b_k y & 1 - a_k b_k y^2 \end{pmatrix}$$



## Taylor method

$$\begin{Bmatrix} q_T \\ p_T \end{Bmatrix} = \begin{pmatrix} T_1^m & T_2^m \\ -T_2^m & T_1^m \end{pmatrix} \begin{Bmatrix} q_0 \\ p_0 \end{Bmatrix} \quad \frac{T\beta}{m} < 0.3$$

## Chebyshev method

$$\begin{Bmatrix} q_C \\ p_C \end{Bmatrix} = \begin{pmatrix} C_1^m & C_2^m \\ -C_2^m & C_1^m \end{pmatrix} \begin{Bmatrix} q_0 \\ p_0 \end{Bmatrix} \quad \frac{T\beta}{m} < 1$$

## Symplectic methods

$$\begin{pmatrix} 1 & 0 \\ -b_k y & 1 \end{pmatrix} \begin{pmatrix} 1 & a_k y \\ 0 & 1 \end{pmatrix} = \begin{pmatrix} 1 & a_k y \\ -b_k y & 1 - a_k b_k y^2 \end{pmatrix}$$

$$K(y) \equiv \prod_{k=1}^m \begin{pmatrix} 1 & a_k y \\ -b_k y & 1 - a_k b_k y^2 \end{pmatrix} = \begin{pmatrix} K_1^{2m-2} & K_2^{2m-1} \\ K_3^{2m-1} & K_4^{2m} \end{pmatrix}$$

$$\begin{Bmatrix} q_S \\ p_S \end{Bmatrix} = \begin{pmatrix} K_1^{2m-2} & K_2^{2m-1} \\ K_3^{2m-1} & K_4^{2m} \end{pmatrix} \begin{Bmatrix} q_0 \\ p_0 \end{Bmatrix} \quad \frac{T\beta}{m} < 2$$

If  $\det K(y) = K_1 K_4 - K_2 K_3 = 1$  then

$$\begin{aligned} & \begin{pmatrix} K_1^{2k-2} & K_2^{2k-1} \\ K_3^{2k-1} & K_4^{2k} \end{pmatrix} \\ = & \begin{pmatrix} 1 & a_k y \\ -b_k y & 1 - a_k b_k y^2 \end{pmatrix} \begin{pmatrix} K_1^{2(k-1)-2} & K_2^{2(k-1)-1} \\ K_3^{2(k-1)-1} & K_4^{2(k-1)} \end{pmatrix} \end{aligned}$$

If the solution exists, it is unique and trivial to obtain.

If  $\det K(y) = K_1 K_4 - K_2 K_3 = 1$  then

$$\begin{aligned} & \begin{pmatrix} K_1^{2k-2} & K_2^{2k-1} \\ K_3^{2k-1} & K_4^{2k} \end{pmatrix} \\ &= \begin{pmatrix} 1 & a_k y \\ -b_k y & 1 - a_k b_k y^2 \end{pmatrix} \begin{pmatrix} K_1^{2(k-1)-2} & K_2^{2(k-1)-1} \\ K_3^{2(k-1)-1} & K_4^{2(k-1)} \end{pmatrix} \end{aligned}$$

If the solution exists, it is unique and trivial to obtain.

In addition

$$\begin{pmatrix} K_1 & K_2 \\ K_3 & K_4 \end{pmatrix} = Q^{-1} \begin{pmatrix} \cos(\phi(y)) & \sin(\phi(y)) \\ -\sin(\phi(y)) & \cos(\phi(y)) \end{pmatrix} Q$$

with  $\phi(y) = \arccos\left(\frac{1}{2}(K_1 + K_4)\right)$

$$K(y)^n \begin{pmatrix} q_0 \\ p_0 \end{pmatrix} = O(n\phi(y)) \begin{pmatrix} q_0 \\ p_0 \end{pmatrix} + E(y) \begin{pmatrix} \sin(n\phi(y))q_0 \\ \sin(n\phi(y))p_0 \end{pmatrix}$$

with

$$E(y) = \begin{pmatrix} \epsilon(y) & \gamma(y) - 1 \\ -\frac{1 + \epsilon(y)^2}{\gamma(y)} + 1 & -\epsilon(y) \end{pmatrix}$$

and

$$\epsilon(y) = \frac{K_1(y) - K_4(y)}{2 \sin(\phi(y))}, \quad \gamma(y) = \frac{K_2(y)}{\sin(\phi(y))}.$$

# Symplectic methods

$$\|u(n\tau) - u_n\| \leq (n\mu(\theta) + \nu(\theta)) \|u_0\|$$

where

$$\mu(\theta) = \sup_{0 \leq y \leq \theta} |\phi(y) - y|, \quad \nu(\theta) = \sup_{0 \leq y \leq \theta} \|E(y)\|.$$




# Symplectic methods

$$\|u(n\tau) - u_n\| \leq (n\mu(\theta) + \nu(\theta)) \|u_0\|$$

where

$$\mu(\theta) = \sup_{0 \leq y \leq \theta} |\phi(y) - y|, \quad \nu(\theta) = \sup_{0 \leq y \leq \theta} \|E(y)\|.$$

$m$	$\theta'$	$\sum_j ( a_j  +  b_j )$	$\mu(\theta)$	$\nu(\theta)$
10	1	4.02	0.00093	0.037
20	1	3.05	0.00061	0.025
30	1	3.19	0.000084	0.037
30	1.4	3.09	0.000051	0.013
30	1	3.04	$2.9 \cdot 10^{-13}$	$2.3 \cdot 10^{-9}$
30	0.75	3.44	$1.2 \cdot 10^{-17}$	$5.9 \cdot 10^{-14}$
30	0.5	3.84	$7.9 \cdot 10^{-24}$	$6.6 \cdot 10^{-18}$
40	1	3.21	$1.1 \cdot 10^{-15}$	$1.0 \cdot 10^{-12}$

-  S. Blanes, F. Casas and A. Murua, On the linear stability of splitting methods, Found. Comp. Math., 8 (2008), pp. 357-393.
-  S. Blanes, F. Casas, and A. Murua, Error analysis of splitting methods for the time dependent Schrodinger equation, SIAM J. Sci. Comput. 33 (2011), pp. 1525-1548
-  S. Blanes, F. Casas, and A. Murua. Work in progress

Group webpage: <http://www.gicas.uji.es>

## Schrödinger Equation with Poschl-Teller Potential

The one-dimensional problem

$$i\frac{\partial}{\partial t}\psi(x, t) = \left( -\frac{1}{2\mu}\frac{\partial^2}{\partial x^2} + V(x) \right) \psi(x, t)$$

with

$$V(x) = -\frac{\alpha^2}{2\mu} \frac{\lambda(\lambda - 1)}{\cosh^2(\alpha x)}$$

$$\mu = 1745, \quad \alpha = 2, \quad \lambda = 24.5,$$

initial conditions

$$\psi(x, 0) = \rho e^{-9^2 x^2}$$



$$t \in [0, T], \quad x \in [-5, 5], \quad \Delta x = 10/N$$

$$E_{min} = V_{min}(x), \quad E_{max} = \frac{1}{2m} \left( \frac{\pi}{\Delta x} \right)^2 + V_{max}(x)$$

$$\beta = \frac{E_{max} - E_{min}}{2}$$

N	64	128	256	512	1024
$\beta$	0.387	0.561	1.25	4.03	15.1

T	$15\pi$	$3\pi$
tol	$10^{-8}$	$10^{-12}$

$$t \in [0, T], \quad x \in [-5, 5], \quad \Delta x = 10/N$$

$$E_{min} = V_{min}(x), \quad E_{max} = \frac{1}{2m} \left( \frac{\pi}{\Delta x} \right)^2 + V_{max}(x)$$

$$\beta = \frac{E_{max} - E_{min}}{2}$$

N	64	128	256	512	1024
$\beta$	0.387	0.561	1.25	4.03	15.1

T	$15\pi$	$3\pi$
tol	$10^{-8}$	$10^{-12}$

	Taylor	Chebyshev	Symplectic
$T \beta = 26.3$ $tol = 10^{-8}$	<b>118</b> $1.5 \cdot 10^{-15}$	<b>50</b> $1.6 \cdot 10^{-11}$	<b>30</b> $1.3 \cdot 10^{-10}$
$T \beta = 37.9$ $tol = 10^{-12}$	<b>208</b> $1.3 \cdot 10^{-13}$	<b>73</b> $2.5 \cdot 10^{-15}$	<b>40</b> $4.7 \cdot 10^{-14}$

# Conclusions

We have shown how to build a class of methods for the Schrödinger equation following the steps previously mentioned:

- 1 We have define mathematically the physical problem
- 2 We have reviewed the [The State of the Art](#) on methods to solve the problem
- 3 We have used our knowledge on the [physical problem](#), [scientific computation](#), [abstract and applied algebra](#), [functional analysis](#), [optimization](#), etc. to improve the existing methods

- 1 To develop a fast and automatic algorithm which finds the optimal coefficients for each particular problem
- 2 To develop new methods when additional information on the problem is known: e.g. If  $\|H^k u_0\| \ll \|H\|^k \cdot \|u_0\|$ .
- 3 To extend the methods to problems with similar structure:  
**Maxwell Equations** or some **linear Hyperbolic PDEs**

# Thank You

## NUMDIFF - 13

Numerical Solution of  
Differential and Differential-  
Algebraic Equations

**10 - 14 September 2012**  
**Halle (Saale), Germany**

